Galaxies at High Redshift

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Contents

Galaxy Formation and Evolution: Recent Progress

R. Ellis

Galaxies at High Redshift

M. Dickinson

High-Redshift Galaxies: The Far-Infrared and Sub-millimeter View

A. Franceschini

Quasar Absorption Lines

Stellar Population Synthesis Models at Low and High Redshift

Elliptical Galaxies

Disk Galaxies

Dark Matter in Disk Galaxies

K. C. Freeman

Galaxy Formation and Evolution: Recent Progress

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In this series of lectures, I review recent observational progress in constraining models of galaxy formation and evolution highlighting the importance advances in addressing questions of the assembly history and origin of the Hubble sequence in the context of modern pictures of structure formation.

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1. Introduction

These are exciting times to be working on any aspect of studies of galaxies at high redshift whether observational or theoretical. Most would agree that the current period represents something of a *golden era* in the subject. Figure 1 shows the increasing extent to which articles concerned with galaxy evolution dominate the published literature over the past 25 years (gauged xenophobically I'm afraid by keyword statistics only in two North American journals).

To try and understand the cause for this prominence in the subject, the dates associated with the commissioning of some major observational facilities have been marked. The progress appears to have been driven largely by new kinds of optical and near-infrared data: faint counts and searches for primaeval galaxies in the late 1970's and early 1980's (Peterson et al. 1979, Tyson & Jarvis 1979, Kron 1980, Koo 1985), faint galaxy redshift surveys made possible by multi-object spectrographs in the late 1980's and early 1990's (Lilly et al. 1995, Ellis et al. 1996, Cowie et al. 1996, Cohen et al. 2000), the launch of Hubble Space Telescope (HST) and its revelation of resolved galaxy images to significant redshifts (Griffiths et al. 1994, Glazebrook et al. 1995, Brinchmann et al. 1998), the remarkable Hubble Deep Field image (Williams et al. 1996) and the plethora of papers that followed (Livio, Fall & Madau 1998) and the arrival of the Keck telescopes bringing a new wave of faint Lyman-break galaxy spectroscopy at unprecedented redshifts (Steidel et al. 1996, Steidel et al. 1999)†.

One often hears claims that a subject undergoing spectacular progress is one that is nearing completion (c.f. Horgan 1997). After all, the rise in Figure 1 clearly cannot continue indefinitely and fairly soon, it could be argued, we will then have solved all of the essential problems in the subject. As if anticipating this, a theoretical colleague gave a recent colloquium at my institute entitled *Galaxy Formation: End of the Road*!

Consider the evidence. Observationally we may soon, via photometric redshifts, have determined the redshift distribution, luminosity evolution and spatial clustering of sources to unprecedented limits. If one accepts photometric redshifts are reliable, the rate of progress in the traditional pursuit of $N(m, color, z)$ is limited solely by the field of view of the telescope and the exposure times adopted. Panchromatic data matching that obtained with optical and near-infrared telescopes from SIRTF, FIRST, and ALMA will

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† A correlation was also made with three key international conferences (Larson & Tinsley 1978, Frenk et al. 1989, Livio, Fall & Madau 1998) but I was horrified to see that these appeared to have had a *negative* effect on the community's output! I assume this arose from a much-needed period of post-conference reflection!

also enable us unravel the cosmic star formation history $\rho_{SFR}(z)$ to unprecedented precision (Madau et al. 1996, Blain et al. 1999). It has already been claimed that the above data, e.g. $N(m, color, z)$ and $\rho_{SFR}(z)$, can be understood in terms of hierarchical models of structure formation where galaxies assemble through the cooling of baryonic gas into merging cold dark matter halos (CDM, Kauffmann et al. 1994, Baugh et al. 1998, Cole et al. 2000a).

The word 'concordance' was recently coined astrophysically in an article reconciling different estimates of the cosmological parameters (Ostriker & Steinhardt 1996). Such concordance in our understanding of galaxy evolution is a natural consequence of semi-analytical theories whose sole purpose is to explain the 'big picture' as realised with the extant galaxy data. In this series of lectures I want to show that we have our work cut out for some considerable time! Exciting progress is definitely being made, but observers must rise to the challenge of testing the fundamentals of contemporary theories such as CDM and theorists must get ready to interpret qualitatively new kinds of data that we can expect in the next decade.
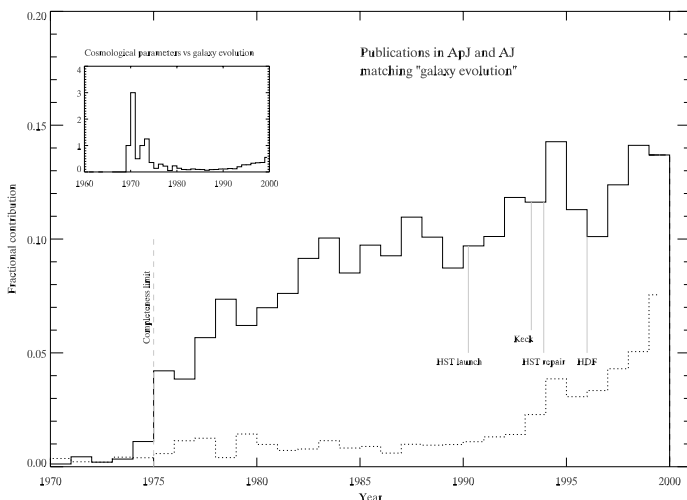
Figure 1: *The remarkably rapid growth in galaxy evolution studies: the fraction of the ApJ and AJ literature containing the key word 'galaxy evolution' over the past 25 years. The inset shows the marked decline in the use of galaxies as probes of the cosmological parameters during 1970-1980 (after Brinchmann, Ph.D. thesis 1999).*

These lectures are intended for interested graduate students or postdocs entering the field. There is an obvious observational flavor although I have tried to keep in perspective an ultimate goal of comparing results with recent CDM predictions. The bias is largely to optical and near-infrared applications; there is insufficient space to do justice to the rapidly-developing contributions being made at sub-millimetre, radio and X-ray wavelengths which other contributors at this winter school will cover in detail.

2. Galaxy Formation and Cosmology

Traditionally faint galaxies were studied in order to constrain the cosmological world model (Sandage 1961); their evolution was considered just one more tedious correction (the so-called *evolutionary correction*) in the path to the Holy Grail of the deceleration parameter q_0 ($\equiv \Omega_M/2$ in Λ=0 Friedmann models). The most useful galaxies in this

respect were giant ellipticals in rich clusters. Tinsley (1976) demonstrated how sensitive the derived q_0 was to the assumed main sequence brightening with look-back time in these populations.

The traditional view for the formation history of an elliptical followed Eggen, Lynden-Bell & Sandage (1962). Monolithic collapse and rapid star formation leads to a subsequent track known as 'passive evolution' (i.e. without further star formation). Tinsley showed that main sequence brightening in such a stellar population is largely governed by the rate at which stars evolve off the main sequence, i.e. the slope $x(\simeq 1)$ of the initial mass function at the typical turnoff mass 0.4–$1M_\odot$. Whence:

$$E(z, t) = d\,M_v/d\ln t \sim 1.3 \,-\, 0.3\,x \qquad (2.1)$$

and, in terms of its bias on q_0:

$$\Delta\,q_0 = 1.4(H_0\,t_o)^{-1}\,d\,M_v/d\ln t = 1.8 \,-\, 0.42\,x \qquad (2.2)$$

Tinsley argued that one would have to know the evolutionary correction to remarkable precision get a secure value of q_0. In fact, noting that the difference in apparent magnitude for a standard candle at $z{=}1$ between an empty and Einstein-de Sitter Universe is only $\simeq 0.5$ mag, the relative importance of cosmology and evolution can be readily gauged.

Despite the above, it is always a mystery to me why several of our most eminent astronomers (Kristian et al. 1978, Gunn & Oke 1975) continued to pursue the Hubble diagram as a cosmological probe using first-ranked cluster galaxies, in some cases for several years after the challenge of resolving the evolutionary correction became known. Tammann (1985) estimated about 400 nights nights of Palomar 200-inch time was consumed by the two competing groups whose resulting values of q_0 fundamentally disagreed. Recently Aragón-Salamanca (1998) showed, in a elegant summary of the situation, how the modern K-band Hubble diagram is most likely complicated further by the fact that first-ranked cluster galaxies are still assembling their stars over the redshift interval $0{<}\,z\,{<}1$, offsetting the main sequence brightening (Figure 2).

In the late 1970's therefore, the motivation for studying faint galaxies became one of understanding their history rather than using them as tracers of the cosmic expansion (see inset panel in Figure 1). This is not to say that uncertainties in the cosmological model do not affect the conclusions drawn. The connection between cosmology and source evolutions remains strong in three respects:

(a) We use our knowledge of stellar evolution to predict the past appearance of stellar populations in galaxies observed at high redshift. However, stellar evolution is baselined in *physical time* (the conventional unit is the Gyr: 10^9yr), whereas we observe distant sources in *redshift* units. The mapping of time and redshift depends on the world model. Broadly speaking there is less time for the necessary changes to occur in a high Ω_M universe and consequently evolutionary trends are much stronger in such models.

(b) Many evolutionary tests depend on the *numbers* of sources, the most familiar being the number-magnitude count which is remarkably sensitive to small changes in source luminosity. However, the relativistic volume element $dV(z)$ depends sensitively on curvature being much larger in open and accelerating Universes than in the Einstein-de Sitter case.

(c) Predictions for the mass assembly history of a galaxy in hierarchical models depend also on the cosmological model in a fairly complex manner since these models jointly satisfy constraints concerned with the normalisation of the mass power spectrum via the

present abundance of clusters (e.g. Baugh et al. 1998). Figure 3 illustrates one aspect of this dependence (Kauffmann & Charlot 1998a); structure grows more rapidly in a dense Universe so the decline with redshift in the abundance of massive spheroidal galaxies, which are thought in this picture to forms via mergers of smaller systems, is much more marked in high density models than in open or accelerating Universes.
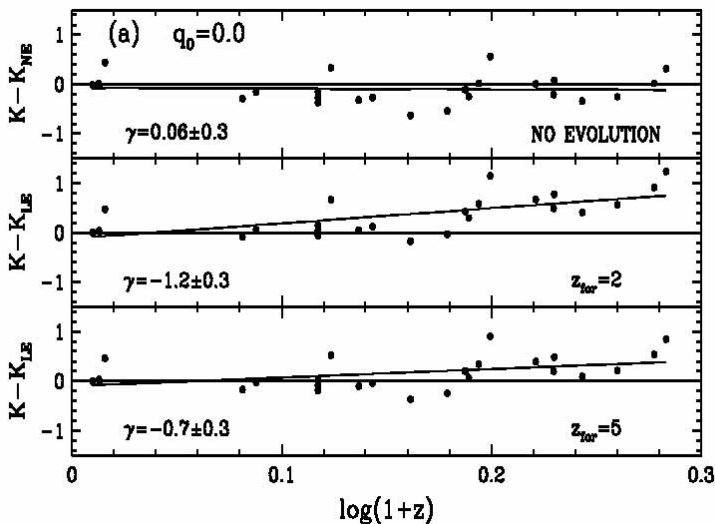
Figure 2: *A recent appraisal of the prospects of securing cosmological constraints from the Hubble diagram of brightest cluster galaxies (Aragón-Salamanca et al. 1998). Luminosity evolution is parameterised as $L = L(0)\,(1 + z)^{\gamma}$. For $q_0 = 0$, the top panel shows residuals and best fit trend applying k-correction and luminosity distance effects only; no luminosity evolution is seen. The middle and bottom panels show the residuals when evolution is modeled for single burst populations formed at $z_F = 2$ and 5, respectively. High z galaxies are* less *luminous than expected, presumably because they are still accreting material. Quantitatively, the effect amounts to a factor of 2-4 less stellar mass depending on the assumed q_0 (c.f. van Dokkum et al. 1999).*

Fortunately, we are making excellent progress in constraining the cosmological parameters from independent methods, the most prominent of which include the angular fluctuation spectrum in the microwave background (de Bernardis et al. 2000, Balbi et al. 2000), the Hubble diagram of distant Type Ia supernovae (Garnavich et al. 1998, Perlmutter et al. 1999), the abundance of rich clusters at various epochs (Bahcall & Fan 1998) and the redshift-space distortion in large redshift surveys such as 2dF (Peacock et al. 2001).

Given it matters, how then should we respond to the widely-accepted *concordance* in the determination of H_0, Ω_M, Λ from various probes (Ostriker & Steinhardt 1996, Bahcall et al. 1999)? The claimed convergence on the value of Hubble's constant (Mould et al. 2000) is not so important for the discussion below since most evolutionary tests are primarily concerned with *relative* comparisons at various look-back times where H_0 cancels. The most bewildering aspect of the concordance picture is the resurrection of a non-zero Λ, the evidence for which comes primarily from the Hubble diagram for Type Ia supernovae.

As a member of the Supernova Cosmology Project (Perlmutter et al. 1999) I obviously

take the supernova results seriously! However, this does not prevent me from being surprised as to the implications of a non-zero Λ. The most astonishing fact is how readily the community has apparently accepted the resurrection of Λ - a term for which there is no satisfactory physical explanation (c.f. Wang et al. 2000). To one poorly-understood component of the cosmic energy density (non-baryonic dark matter), we seem to have added another (vacuum energy). It seems a remarkable coincidence that all three significant constituents $(\Omega_B, \Omega_{DM}, \Omega_\Lambda)$ are comparable in magnitude to within a factor of 10, and hardly a step forward that only one is physically understood!
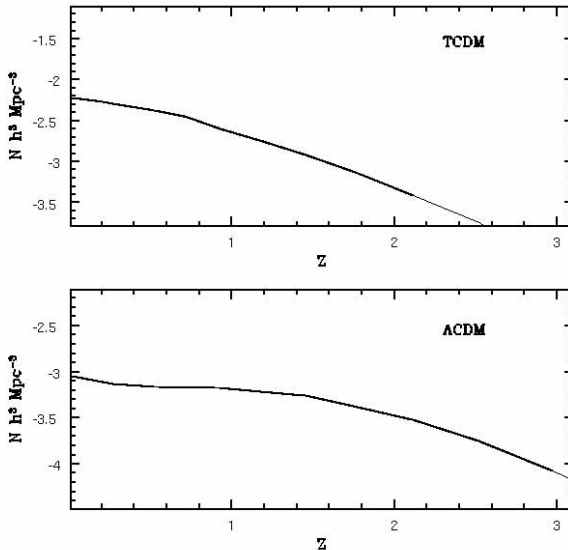
Figure 3: *The abundance of massive $(>10^{11}M_\odot)$ systems as a function of redshift in two hierarchical models (Kauffmann & Charlot 1998a) showing the strong decline in a high density (τCDM) model c.f. that in a low density accelerating model (ΛCDM).*

The lesson I think we should draw from the *cosmic concordance* is similar to the comment I made in §1 when we discussed some theorists' triumphant reconciliation of their theories with faint galaxy data (a point we will debate in detail in §3). In both cases, the hypothesis certainly reproduces a wide range of observations but note it takes, as input, parameters for which there is not yet a clear physical model. One should not, therefore, regard a concordant picture as anything other than one of many possible working hypotheses. In the case of the cosmological models, we need to invest effort into understanding the physical nature of dark matter and vacuum energy. In the case of galaxy evolution our goal should be to test the basic ingredients of hierarchical galaxy formation.

3. Star Formation Histories

One of the most active areas of relevance to understanding the rate at which galaxies assemble is concerned with determining the cosmic star formation history. The idea is simple enough. A systematic survey is conducted according to some property that is sensitive to the on-going rate of star formation. The volume-average luminosity density is converted into its equivalent star formation rate averaged per unit co-moving volume

and the procedure repeated as a function of redshift to give the cosmic star formation history $\rho_*(z)$. In this section we will explore the uncertainties and also the significance of this considerable area of current activity in terms of the constraints they provide on theories of galaxy formation.
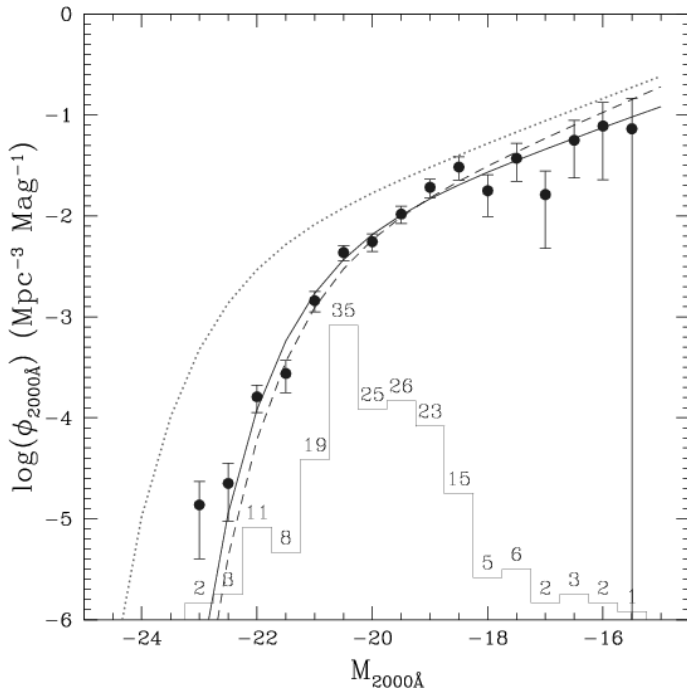
Figure 4: *The luminosity function for galaxies selected at 2000 Å from the recent survey of Sullivan et al (2000). The histogram and associated numbers indicate the absolute magnitude distribution observed which is corrected by volume and k-correction effects to give the data points. The dotted curve illustrates the considerable effect of extinction as gauged by Balmer decrements determined individually for those galaxies with emission lines. Such uncertainties translate in factors of two uncertainty in the local UV luminosity density.*

The joint distribution of luminosity L and redshift z, $N(L, z)$, for a flux-limited sample permits the construction of the luminosity function $\Phi(L)$ according to procedures which are reviewed by Efstathiou, Ellis & Peterson (1988) and compared by Ellis (1997) . The luminosity function is often characterised according to the form defined by Schechter (1976), viz:

$$\Phi(L)\,dL/L^* = \Phi^*\,(L/L^*)^{-\alpha}\,\exp(-L/L^*)\,dL/L^* \qquad (3.3)$$

in which case the integrated number of galaxies per unit volume N and the luminosity density ρ_L then become:

$$N = \int \Phi(L)\,dL = \Phi^*\,\Gamma(\alpha + 1) \qquad (3.4)$$

and

$$\rho_L = \int \Phi(L)\,L\,dL = \Phi^*\,L^*\,\Gamma(\alpha + 2) \qquad (3.5)$$

and the source counts in the non-relativistic case, applicable to local catalogs, is:

$$N(<m) \; \propto \; d^{*\,3}(m) \int dL \, \Phi(L) \, (L/L^*)^{\frac{3}{2}} \; \propto \; \Phi^* \, L^{*\,\frac{3}{2}} \, \Gamma(\alpha + \frac{5}{2}) \qquad (3.6)$$

Frequently-used measures of star formation in galaxies over a range of redshift include rest-frame ultraviolet and blue broad-band luminosities (Lilly et al. 1995, Steidel et al. 1996, Sullivan et al. 2000), nebular emission lines such as Hα (Gallego et al. 1995, Tresse & Maddox 1998, Glazebrook et al. 1999), thermal far-infrared emission from dust clouds (Rowan-Robinson et al. 1997, Blain et al. 1999) and, most recently, radio continuum emission (Mobasher et al. 1999).

Since only a limited range of the luminosity function centered on L^* is reliably probed in flux-limited samples, a key issue is how well the integrated luminosity density ρ_L can be determined from such surveys. In the Schechter formalism, equations [3.4] and [3.5] show that whilst N would diverge for $\alpha < -1$, the luminosity density is convergent unless $\alpha < -2$.

Figure 4 shows the local rest-frame ultraviolet (2000 Å) luminosity function from Sullivan et al. (2000) whose faint end slope $\alpha = -1.6$ is markedly steeper than that found for samples selected in the near-infrared (Mobasher et al. 1993, Gardner et al. 1997, Cole et al. 2001) (where $\alpha \simeq -1$). This contrast in the luminosity distribution of young and old stellar populations is an important result which emphasizes the relatively weak connection between stellar mass and light and implies there may be significant uncertainties in the estimation of integrated luminosity densities for star-forming populations.
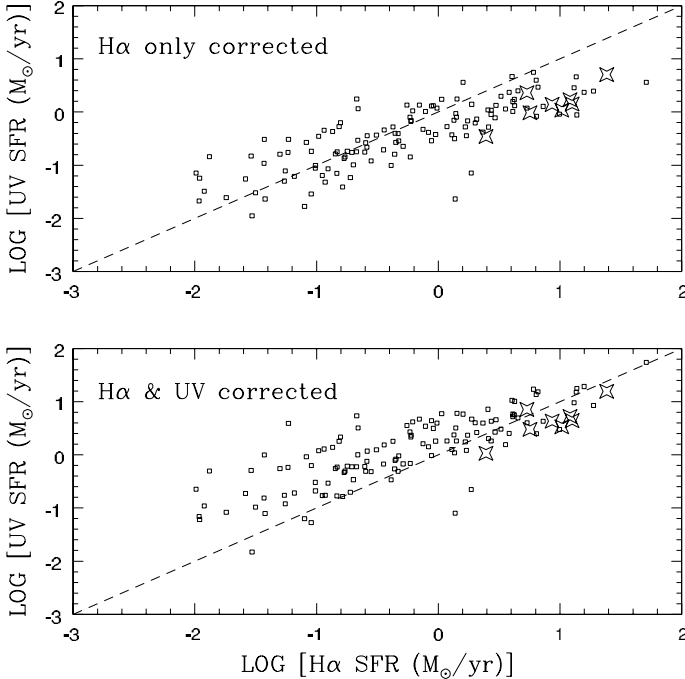
Figure 5: *Star formation rates derived from UV (2000 Å) continua versus those derived from Hα fluxes from the local survey of Sullivan et al. (2000, open squares) and the $z \simeq 1$ samples of Glazebrook et al. (1999, large stars). For the Sullivan et al. sample, extinction corrections were derived from individual Balmer decrements assuming Case B recombination and applied to the Hα fluxes in the upper panel and both estimates in the lower panel.*

Kennicutt (1998) carefully reviewed the relationships between the various observational diagnostics listed above and the star formation rate. Clearly a major uncertainty in any transformation based on the ultraviolet/optical continuum or nebular emission line measures is the likely presence of absorbing dust (Figure 4). Other uncertainties include the form of the initial stellar mass function and the nature of the star formation history itself.

Some of these uncertainties are quite imponderable and the only way to estimate their effect in typical populations is to undertake a comparison of the various diagnostics for the same sample. Sullivan et al. (2000) compared UV and Hα-based estimators for their local balloon-based UV-selected sample and Glazebrook et al. (1999) undertook a similar comparison for a restricted incomplete sample of high redshift galaxies (drawn from a I-selected sample). Bell & Kennicutt (2001) independently examined some of Sullivan et al's conclusions based on a smaller local sample with satellite UV fluxes. The comparison analysed by Sullivan et al. is shown in Figure 5. Although an overall linear relation is observed the scatter is quite considerable, greater than accountable from observational errors. The uncertainties would appear to be alarming in view of the fairly modest trends claimed in $\rho_{SFR}(z)$ (see below).

In addition to the scatter arising from extinction (accounted for via individual Balmer emission line decrements), Sullivan et al suggest that some fraction of their UV-selected population must be suffering star formation which is erratic in its time history. In such a situation, different diagnostics will be sensitive to bursts of activity for different periods, corresponding to the time over which the contributing stars remain on the main sequence. Hα flux arises from recombination photons linked to those emitted below the Lyman limit from main sequence stars with lifetimes $\simeq 10^6$ years. The UV and blue continua persist for much longer periods ($\simeq 10^8 - 10^9$ years).

Depending upon how widespread star formation histories of this kind may be, two forms of error may arise in estimating cosmic star formation histories. Firstly, the star formation rate derived for an individual galaxy will be a past time average, smoothing over any erratic behavior, rather than a true instantaneous value. More importantly however, particularly at high redshift, galaxies may be preferentially selected only if their star formation history is erratic, for example in Hα surveys where some threshold of detectability may seriously restrict the samples.

Figure 6 shows a recent estimate of the cosmic star formation history drawn from various surveys (Blain 2001). There appears to be a marked increase in activity over $0 < z < 1$ with a possible decline beyond $z > 2$. Although, inevitably perhaps, attention has focused on the case for the high redshift decline, even the strong rise to $z \simeq 1$ remains controversial. Originally proposed independently by Lilly et al. (1995) and Fall et al. (1995), revised estimates for the local luminosity density (Sullivan et al. 2000) and independent surveys (Cowie et al. 1999) have challenged the rapidity of this rise. Part of the problem is that no single survey permits a self-consistent measurement of ρ_{SFR} over more than a very limited range in z. Most likely, therefore, much of the scatter in Figure 6 is simply a manefestation of the kinds of uncertainties discussed above in the context of Sullivan et al.'s survey.

Beyond $z \simeq 2$, the available star formation rates have been derived almost exclusively from UV continua in Lyman break galaxies selected by their 'dropout' signatures in various photometric bands (Madau et al. 1996, Steidel et al. 1996, Steidel et al. 1999) and from currently scant datasets of sub-mm sources interpreted assuming thermal emission from dust heated by young stars (Blain et al. 1999, Barger et al. 1999b). There has been much discussion on the possible disparity between the estimates derived from these two diagnostics (which other lecturers will address). Two points can be made: firstly, the

measured UV luminosity densities will clearly underestimate the true values given likely extinctions. Secondly, the sample of sub-mm sources with reliable redshifts remains quite inadequate for luminosity density estimates in the sense described above. Most of the constraints arise from modelling their likely properties in a manner consistent with their source counts and the integrated far-infrared background.

Have we become over-obsessed with determining the cosmic star formation history? Observers are eager to place their survey points alongside others on the overall curve and different groups defend their methods against those whose data points disagree. We should consider carefully what role this cosmic star formation history plays in understanding how galaxies form.
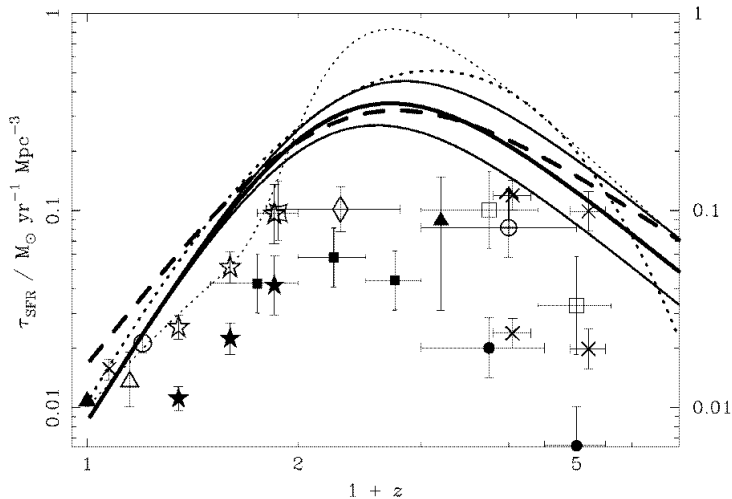

Figure 6: *The history of recent star formation from the recent compilation of Blain (2001). Data points are taken from a variety of sources referenced in that article. Thick solid and dashed lines represent trends expected from simple luminosity evolution and hierarchical models, respectively. It is clear there is considerable observational scatter at all redshifts, not just beyond $z \simeq 1$ as often assumed.*

Clearly, the prime conclusion we can draw from Figure 6 is that the stars which make the galaxies we see today formed continuously over a very wide redshift range. This may seem such an obvious deduction that it hardly merits stating but it is important to stress the absence of any obvious detectable 'epoch of star formation' as was once imagined (Eggen, Lynden-Bell & Sandage 1962, Frenk et al. 1989). Hierarchical modelers were quick to point out (e.g. Baugh et al. 1998) that they predicted extended star formation histories as early as 1990 (White & Frenk 1991). It is certainly true that a continuous assembly of galaxies is a major feature of these models and thus one supported by the data.

However, what about the *quantitative* form of Figure 6 which remains so difficult to pin down: does the shape of the curve really matter? Firstly, we should recognise that the luminosity density integrates over much detailed astrophysics that may be important. A particular ρ_{SFR} at a given redshift could be consistent *either* with a population of established massive sources undergoing modest continuous star formation *or* a steep luminosity function where most of the activity is in newly-formed dwarf galaxies. In terms of structure formation theories, these are very different physical situations yet that distinction is lost in Figure 6.

Secondly, theoretically, the cosmic star formation history is not particularly closely related to how galaxies assemble. It is more sensitive to the rate at which gas cools into the assembling dark matter halos, a process of considerable interest but which involves a myriad of uncertain astrophysical processes (Figure 7) which are fairly detached from the underlying physical basis of say the hierarchical picture. In support of this, we should note that Baugh et al. (1999) were able, within the same Λ-dominated CDM model, to 'refine' their earlier prediction to match new high redshift datapoints revealing a much less marked decline beyond $z \simeq 2$.
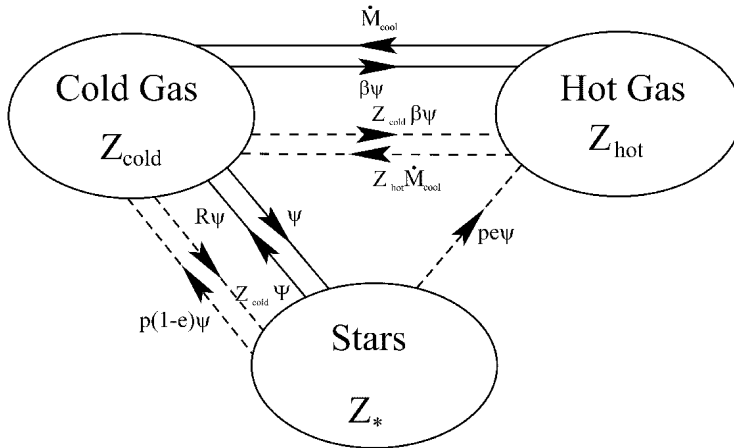

Figure 7:   *An illustration of the complex physical processes governing the star formation rate of a young galaxy (courtesy of Carlos Frenk). Star formation is governed by the rate at which baryonic gas cools and falls into dark matter halos and this is inhibited by heating, e.g. from supernovae. The precise form of the cosmic star formation history gives us more insight into the interplay between these processes, integrated over all star-forming galaxies, than in distinguishing between various forms of structure formation (e.g. hierarchical vs. monolithic).*

4. Morphological Data from HST

As we discussed in §1, one of the most exciting new datasets that arrived in the mid-1990's was the first set of resolved images of galaxies at significant look-back times from HST. Much of the early work was conducted in rich clusters (Couch et al. 1994, Dressler et al. 1994, Couch et al. 1998, Dressler et al. 1998) where the well-known 'Butcher-Oemler' effect (Butcher & Oemler 1978) - a surprisingly recent increase in the fraction of blue cluster members - was found to be due to a dramatic shift in the morphology-density relation (Figure 8). As recently as $z \simeq 0.3 - 0.4$ ($3 - 4$ Gyr ago), cluster S0s were noticeably fewer in proportion, their place apparently taken by spirals, many of which showed signs of recent disturbances, such a distorted arms and tidal tails.

The physical origin of this transformation from spirals to S0s remains unclear and is currently being explored by detailed spectroscopy of representative cluster members (Barger et al. 1996, Abraham et al. 1996, Poggianti et al. 1999). A key diagnostic here is the interplay between the changing morphologies, the presence of nebular emission lines (such as [O II] 3727 Å – Hα is generally redshifted out of the accessible range) and
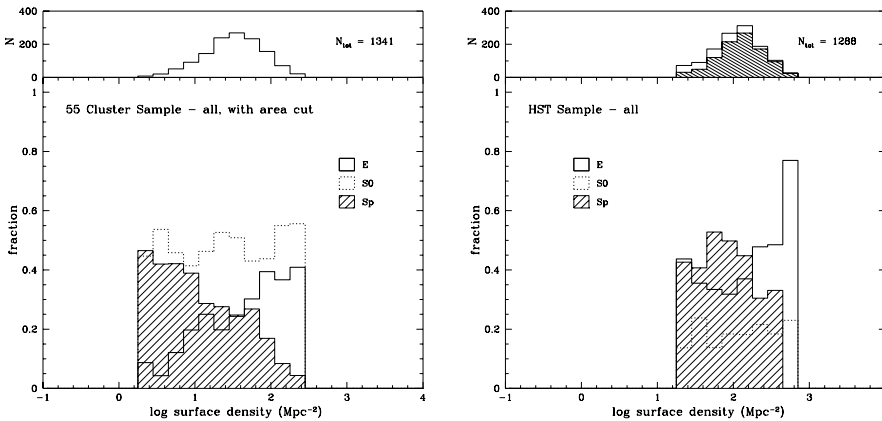
Figure 8: *Evolution in the morphology density relation from the 'Morphs' survey of Dressler et al. (1998). (Left) The fraction of E/S0/Spirals as a function of projected galaxian surface density for Dressler's 55 local cluster sample. (Right) As left, for all distant HST clusters with $0.3 < z < 0.55$. The comparison refers to the same cluster core radius ($< 0.6Mpc$, $h = 0.5$) and includes galaxies to the same rest-frame V luminosity ($M_V = -20.0$). Note the dramatic decline in the S0 population and the marked increase in the spiral fraction for environments with high projected density.*

Balmer absorption lines (such as Hδ 4101 Å). The latter lines are prominent in main sequence A stars which linger for $\simeq$1 Gyr after any enhanced starburst activity.

Barger et al. (1996) proposed a simple cycle where an unsuspecting galaxy undergoes some perturbation, perhaps due to a merger or its first encounter with the intracluster gas, subsequently becomes morphologically-distorted and spectrally-active before subsiding to a regular spheroidal with a decaying Balmer absorption line. Whereas such a cycle can explain the *proportion* of unusual objects, it has difficulty matching their *luminosities*. A galaxy should be rendered more luminous during a burst and thus blue examples cannot easily be the precursors of the equally-luminous red *post-burst* cases. A controversy has since arisen over the fractions of objects seen in the various spectrally-active classes (Balogh et al. 1999) suggesting much work is needed in this area, both in quantifying cluster-cluster variations and also radial variations in the responsible processes.

Although the cluster work discussed above represents something of a digression in our overall theme, the realisation that galaxies can so easily be transformed morphologically has profound implications for our understanding of galaxy formation. Much of the early work explaining the Hubble sequence (Tinsley 1977) assumed galaxies evolve as isolated systems, however the abundance of morphologically-peculiar and interacting galaxies in early HST images (Griffiths et al. 1994) has been used to emphasize the important role that galaxy mergers must play in shaping the present Hubble sequence (Toomre & Toomre 1972, Barnes & Hernquist 1992). Merger-induced transformations of this kind are a natural consequence of hierarchical models (Baugh, Cole & Frenk 1996). Early disk systems are prone to merge during epochs when the cosmic density is high and the peculiar velocity field is cold, forming bulge-dominated and spheroidal systems which may then later accrete disks.

The possibility that galaxies transform from one class to another is a hard hypothesis to verify observationally since, as we have seen, traditionally observers have searched for redshift-dependent trends with subsets of the population chosen via an observed

property (color, morphology, spectral characteristics) which could be transient. More-over, experience ought to teach us that the outcome of tests of galaxy formation rarely come down simply to either Theory A *or* Theory B; usually it is some complicated mixture or the question was naive in the first place! Fortunately, the late formation of massive regular galaxies in the hierarchical picture (Figure 9) seems a particularly robust prediction and one in stark contrast to the classical 'monolithic collapse' picture (Tinsley 1977, Sandage 1986). The distinction is greatest for ellipticals presumed to form at high redshift with minimum dissipation (their central density reflecting that of the epoch of formation). Studying the evolutionary history of massive ellipticals is thus an obvious place to start.
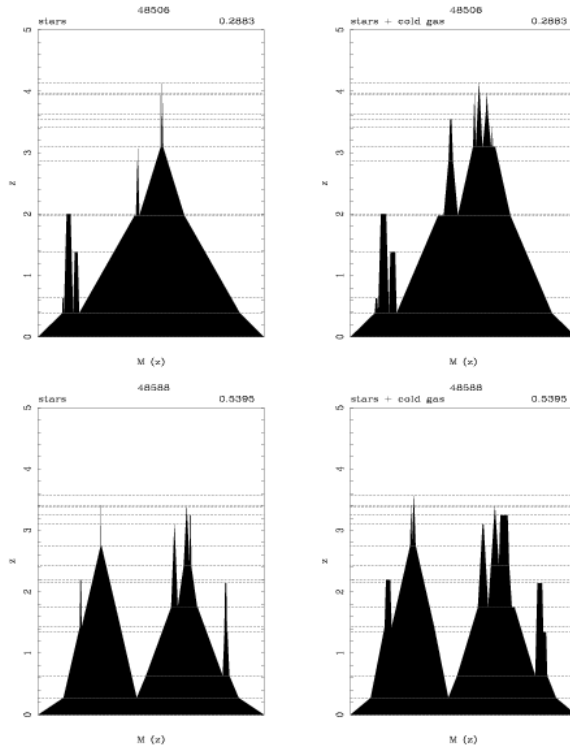

Figure 9: *The important role of late merging in a typical CDM semi-analytical model (Baugh et al. 1996). The panels show the redshift-dependent growth for the stellar mass (left) and that of all baryonic material (right), as indicated by the thickness of the black area at a given epoch, for two present-day massive galaxies. The top system grows grad-ually and is thought to represent a present-day spiral. The bottom system suffers a late equal-mass merger thought to produce a present-day elliptical. Note the remarkably late assembly; most of stellar mass in both cases assembles in the interval $0 < z < 1$.*

An oft-quoted result in support of old ellipticals is the remarkable homogeneity of their optical colors (Sandage & Visvanathan 1978, Bower et al. 1992). The idea is simple: the intrinsic population scatter in a color sensitive to recent star formation, such as $U - B$, places a constraint either on how synchronous the previous star formation history must have been across the population or, if galaxies form independently, the mean age of their stellar populations. By combining cluster data at low redshift (Bower et al. 1992) with

HST-selected samples at intermediate redshift (Ellis et al. 1998), the bulk of the cluster elliptical population was deduced to have formed its stars before $z \simeq 2$, in apparent conflict with hierarchical models. Similar conclusions have been drawn from evolution of the mass/light ratio deduced from the fundamental plane (Ziegler & Bender 1997, van Dokkum et al. 1998).

Unfortunately, one cannot generalize from the results found in distant clusters. In hierarchical models, clusters represent early peaks in the density fluctuations and thus evolution is likely accelerated in these environments (Kauffmann 1995) plus, of course, there may be processes peculiar to these environments involving the intracluster gas. It is also important to distinguish between the history of *mass* assembly and that of the *stars*. Recent evidence for widespread merging of ellipticals in clusters (van Dokkum et al. 1999) lends support to the idea that the stars in dense regions were formed at high redshift, in lower mass systems which later merged.

For these reasons, attention has recently switched to tracking the evolution of *field* ellipticals. The term *field elliptical* is something of a misnomer here since a high fraction of ellipticals actually reside in clusters. What is really meant in this case is that we prefer to select ellipticals systematically in flux-limited samples rather than concentrate on those found in the cores of dense clusters†.

The study of evolution in field ellipticals is currently very active and I cannot possibly do justice, in the space available, to the many complex issues being discussed. Instead let me summarise what I think are the most interesting results.
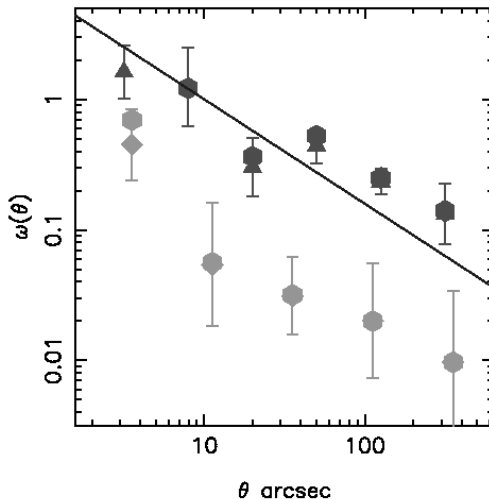
Figure 10: *Evidence for a clustered population of red objects in a 1000 arcmin² area of the ongoing Las Campanas Infrared Survey (McCarthy et al. 2000). The angular correlation function for objects with $I - H > 3.5$ (top set of data points) is substantially above that for all H-selected galaxies (bottom set) and consistent with a high fraction of the red objects being clustered ellipticals with $\overline{z} \simeq 1.0 - 1.5$.*

• Searches for a population of faint intrinsically red objects, representing the expected $z > 1$ precursors of passively-evolving ellipticals which formed their stars at high redshift have been conducted both with and without HST morphological data (Zepf 1997, Barger et al. 1999a, Menanteau et al. 1999, Daddi et al. 2000, McCarthy et al. 2000).

† A major concern in all the work relating to the evolution of galaxies in clusters is precisely how the clusters were located.

However, only recently have substantial areas of sky been mapped. This is because such searches are most sensitive to high z sources when conducted using optical-near infrared colors and access to large infrared arrays is a recent technical development. Both Daddi et al (2000) and McCarthy et al. (2000) (Figure 10) claim strong angular clustering in their faint red populations and there is limited evidence that the abundance is consistent with a constant comoving number density, in contrast to the hierarchical predictions. However, without confirmatory spectroscopy neither the redshift range nor the nature of these red sources is yet clear. Even if it later emerges, as was claimed on far less convincing data (Zepf 1997, Barger et al. 1999a, Menanteau et al. 1999), that there is a shortage of intrinsically red objects beyond $z \simeq 1$, only a modest amount of residual star formation is needed to substantially bluen a well-established old galaxy (Jimenez et al. 1999). Even when redshift data is secured, color alone may be an unreliable way to track a specific population.
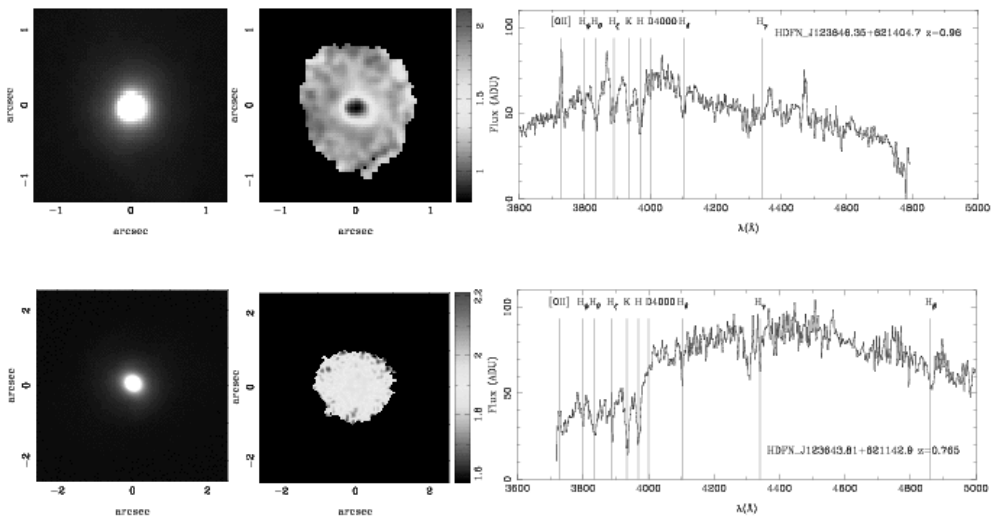

Figure 11: *Color inhomogeneities in HDF field ellipticals suggest continued star formation is occurring, possibly as a result of hierarchical assembly. Each row displays the I HST image, a $V - I$ color image and the Keck LRIS spectrum. The top set refers to a z=0.92 elliptical with a blue core; its spectrum shows features indicative of active star formation ([O II] emission and deep Balmer absorption lines). The bottom set refers to a quiescent example at z=0.966 whose spectrum is consistent with an old stellar population. The amount of blue light can be combined with the depth of the spectral features to statistically estimate the amount and timescale of recent star formation.*

• At brighter magnitudes, systematic redshift surveys with associated HST data give constraints on the luminosity function and colors of morphologically-selected ellipticals (Brinchmann et al. 1998, Schade et al. 1999, Menanteau et al. 1999, Im et al. 2001). Unfortunately, because of the disparity in field of view betweeen WFPC-2 and the ground-based multislit spectrographs, the samples remain small and hence the conclusions are subject to significant field-to-field clustering uncertainties. However, no substantial decline in the volume density of ellipticals is yet observed to $z \simeq 1$, although there is some dispute as to the fraction which may deviate in color from the passive track (Im et al. 2001 c.f. Menanteau et al. 1999). Current spectroscopic surveys may not

be quite deep enough to critically test the expected evolution in the hierarchical models, particularly if $\Lambda \neq 0$.

• A completely independent method of determining whether field ellipticals form continuously as expected in hierarchical models is possible in the Hubble Deep Fields (Menanteau et al. 2001). Here, the imaging signal/noise is sufficient to permit an examination the *internal* colors of ellipticals with $I < 24$, a subset of which have redshifts. Menanteau et al. (2001) find about 25% of the HDF ellipticals show blue cores and other color inhomogeneities suggestive of recent star formation, perhaps as a result of the merger with a gas-rich low mass galaxy. Keck spectroscopy (Ellis et al. 2001b) supports this suggestion: galaxies with blue cores generally show emission and absorption line features indicative of star formation (Figure 11). The amount of blue light seen in the affected ellipticals can be used to quantify the *amount* of recent star formation and the associated spectrum can be used to estimate the *timescale* of activity through diagnostic features of main sequence stars. Only modest accretion rates of $\simeq 10\%$ by mass over $\simeq 1$ Gyr are implied, albeit for a significant fraction of the population. This continued growth, whilst modest c.f. expectations of hierarchical models, is noticeably *less* prominent in rich clusters (Menanteau et al. 2001).

What evolution is found in the properties of other kinds of galaxies? Brinchmann et al. (1998) secured HST images for a sizeable and statistically-complete subset of CFRS and LDSS redshift survey galaxies and found the abundance of spirals to $I = 22$ – a flux limit which samples $0.3 < z < 0.8$ – is comparable to that expected on the basis of their local abundance if their disks were somewhat brighter and bluer in the past as evidenced from surface photometry (Lilly et al. 1999). In practice, however, the detectability of spiral disks is affected by a number of possible selection effects (Simard et al. 1999, Bouwens & Silk 2000) and it may be some time before a self-consistent picture emerges.

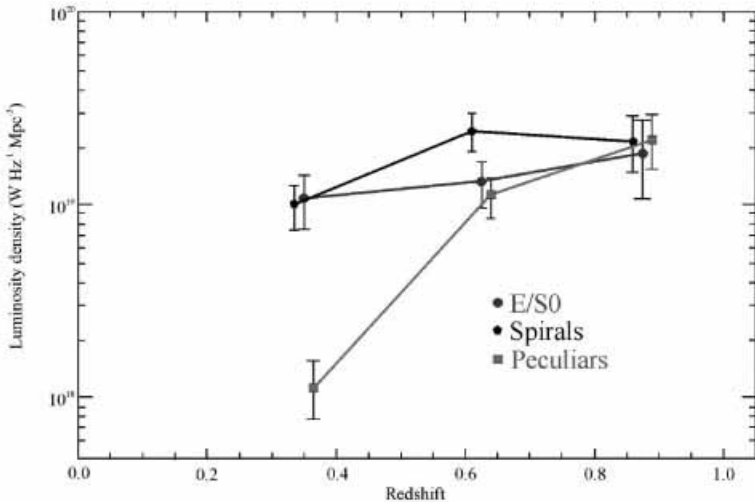**Figure 12:** *The morphological dependence of the blue luminosity density from that subset of the CFRS/LDSS redshift survey imaged with HST (Brinchmann et al. 1998). The marked decline in the luminosity density of galaxies with peculiar morphology over $0 < z < 1$ is the primary cause for steep slope in the blue faint galaxy counts.*

A less controversial result from Brinchmann et al. (1998) claimed in earlier analyses
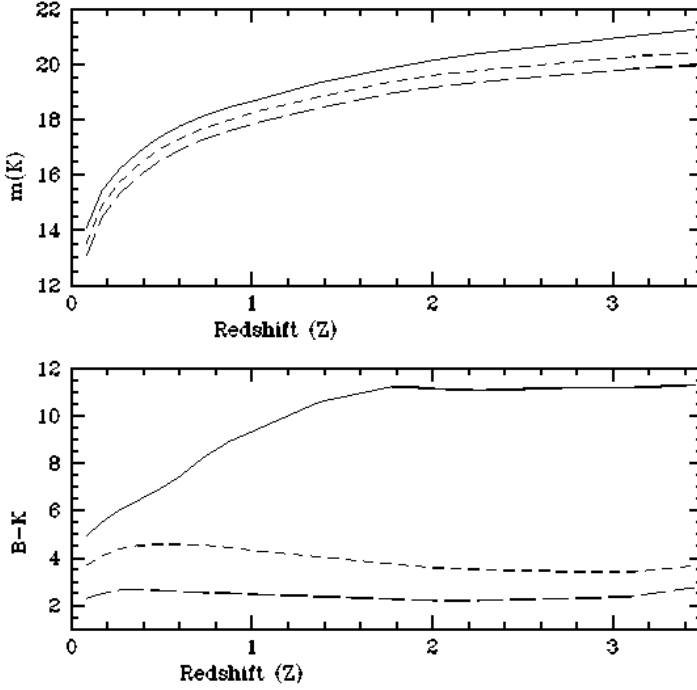
Figure 13: *The K-band luminosity is a good measure of the underlying stellar mass irrespective the past star formation history (Kauffmann & Charlot 1998b). The curves show the observed K magnitude as a function of the redshift at which such an object is selected for a system containing $10^{11} M_\odot$ produced according to a variety of star formation histories. Even across extreme cases (single burst at $z = \infty$, solid line, to a constant star formation rate to the epoch of observation, short-dash), the K-band output remains the same to within a factor of $\simeq 2$. The lower panel shows, how different the observe optical-infrared color would be in these cases.*

without redshift data (Glazebrook et al. 1995, Driver et al. 1995) is the remarkably high abundance of morphologically-peculiar galaxies in faint HST data. Brinchmann et al. quantified this in terms of the luminosity density arguing that a substantial fraction of the claimed decline in the blue luminosity density since $z \simeq 1$ (c.f. Figure 6) arises from the demise of this population (Figure 12).

Given our earlier concerns with over-interpreting the cosmic star formation history, should we be cautious in drawing conclusions from Figure 12? Although Brinchmann et al's redshift sample is small, the basic result is consistent with the HST morphological number counts where much larger samples are involved. Whereas early skeptics argued that morphologically-peculiar galaxies represent regular systems viewed at unfamiliar ultraviolet wavelengths, recent NICMOS imaging (see Dickinson's lectures) suggests such 'morphological bandshifting' is only of minor consequence. In quantitative detail, as before, uncertain corrections must be made for the effects of the flux limited sample and of course extinction is a major uncertainty. However, it seems inescapable that the bulk of the decline in blue light (the so-called *faint blue galaxy problem*, Ellis 1997) arises

from the demise of a population of late-type and morphologically-peculiar systems. A key question therefore is what happened to this population? We will address this problem in the next section.

5. Constraining the Masses of Distant Galaxies

A recurring issue arises from the discussions in the earlier sections. Whilst observers are, with some restrictions, able to measure distant galaxy properties such as rest-frame colors, luminosities and star formation rates, these may be poor indications of the underlying stellar and total masses predicted most straightforwardly by contemporary models of structure formation. Either we put our faith in the forward modelling of the readily-available observables (i.e. we invest a lot of effort in understanding the complexities of feedback, Figure 7), or we consider how to measure galactic masses.

Ideally we seek methods for determining the *total* mass (baryonic plus the dark matter halo) but this seems out of reach for the moment except for local systems with tracers of the larger halo in which galaxies are thought to reside. Useful tracers here include the dynamical properties of attendant dwarf galaxies (Zaritsky et al. 1998) and globular clusters. A promising route in the future might be galaxy-galaxy gravitational lensing (Blandford & Narayan 1992). Here a foreground population is restricted in its selection, perhaps according to morphology or redshift, and the statistical image distortions in a background population analysed. Early results were based on HST data, for cluster spheroidals (Natarajan et al. 1999) and various field populations (Griffiths et al. 1996), however with extensive panoramic data from the Sloan Digital Sky Survey, convincing signals can be seen with ground-based photometry (Fischer et al. 2000). Again, photometric redshifts will be helpful in refining the sample selection and in determining the precise redshift distribution essential for accurate measures on an absolute scale.

Unfortunately, promising though the technique appears, the restrictions of galaxy-galaxy lensing are numerous. It only gives mass estimates for statistical samples: the signal is too weak to be detected in individual cases, unless a strong lensing feature is seen (Hogg et al. 1996). Moreover, the redshift range and physical scale on which the mass is determined is defined entirely by geometrical factors and, ultimately, one may never be able to apply the method to galaxies beyond $z \simeq 1$.

Extensive dynamical data is becoming available for restricted classes of high redshift galaxy, via linewidth measures (Koo et al. 1995), resolved rotation curves (Vogt et al. 1997) for sources with detectable [O II] 3727 Å emission, and via internal stellar velocity dispersions for absorption line galaxies such as spheroidals. Under certain assumptions, these give mass estimates and have enabled the construction of the fundamental plane for distant spheroidals (Treu et al. 2000) and the Tully-Fisher relation for high redshift disk galaxies (Vogt et al. 1997). The greatest progress in the former has been in constructing the fundamental plane in rich clusters (van Dokkum et al. 1998) where slow evolution in the inferred mass/light ratio for cluster ellipticals is consistent with a high redshift of formation (see §5). For the emission line studies it is not straightforward to convert data obtained over a limited spatial extent into reliable masses even for regular well-ordered systems. For compact and irregular sources, the required emission lines may come from unrepresentative components yielding poor mass estimates (Lehnert & Heckman 1996).

The prospects improve significantly if we drop the requirement to measure the *total* mass and are willing to consider only the *stellar mass*. In this case, the near-infrared luminosity is of particular importance. Broadhurst et al. (1992) and Kauffmann & Charlot (1998b) have demonstrated that the K (2.2μm) luminosity is a good measure of
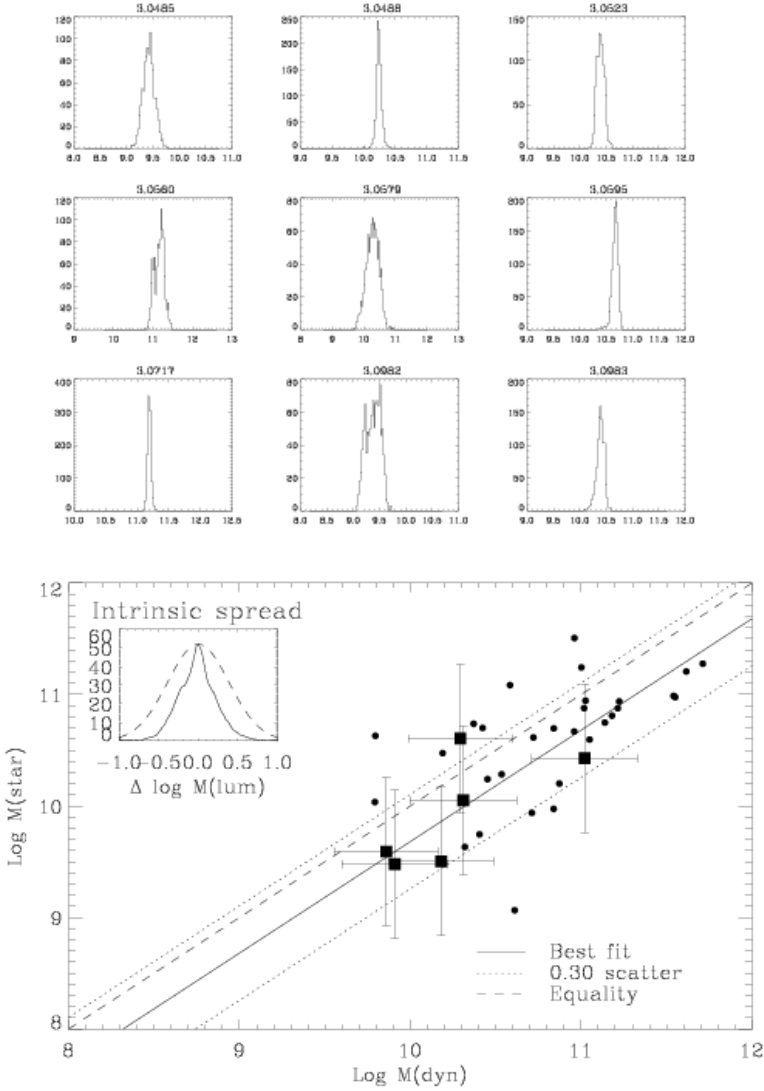
Figure 14: *The infrared method for determining the stellar mass of a distant galaxy (Ellis et al. 2001a). The technique fits the observed SED for a galaxy of known redshift in the context of evolutionary synthesis models where the stellar mass is the fitted variable. (Top) Likelihood functions for the derived logarithmic stellar mass for sample galaxies in the CFHT/LDSS redshift survey (Brinchmann, Ph.D. thesis 1999); a typical uncertainty of 30-50% is secured at I≃22. (Bottom) Correlations of stellar and dynamical mass for both low z (circles) and high z (squares with error bars) galaxies from the analysis of Brinchmann & Ellis (2000).*

its underlying stellar mass *regardless of how that mass assembled itself* (Figure 13). This remarkable fact arises because K-band light in all stellar populations (whether induced in bursts or continuous periods of activity) arises from long-lived giants whose collective output mirrors the *amount* of past activity, smoothing over its production timetable.

A deep K-band redshift survey thus probes the very existence of massive systems at early times. A slightly incomplete survey to $K=20$ (Cowie et al. 1996) and a complete

photometric survey to $K=21$ (Fontana et al. 1999) indicates an apparently shortfall of luminous K objects beyond $z \simeq 1 - 1.5$ c.f. pure luminosity evolution models. Unfortunately, small sample sizes, field-to-field clustering, spectroscopic incompleteness and untested photometric redshift techniques beyond z≃1 each weaken this potentially important conclusion. An important goal in the immediate future must be to reconcile these claims with the apparently abundant (and hence conflicting) population of optical-infrared red objects to $K \simeq 19 - 20$ (Daddi et al. 2000, McCarthy et al. 2000).
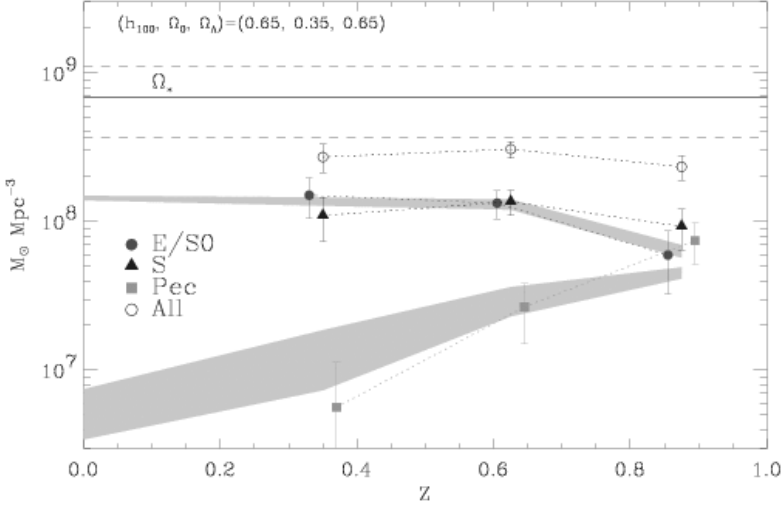

Figure 15: *Evolution of the stellar mass density $\rho_{stars}(z, T)$ from the analysis of Brinchmann & Ellis (2000). A remarkable decline with time in stellar mass density is seen for the morphologically-peculiar class which argues against a truncation of their star formation activity as the primary cause for their demise. Brinchmann & Ellis argue that this population must be transforming, possibly via mergers, into the regular classes. A simple model which implements a likely redshift-dependent merger rate (LeFevre et al. 2000) with elliptical products can broadly reproduce the trends observed (shaded area of the plot).*

The precision of the technique introduced by Kauffmann & Charlot (1998b) can be improved if the optical-infrared color is available as an extra parameter (Ellis et al. 2001a). In this way a first-order correction can be made for the past star formation history and hence the effect of the spread in the lower panel of Figure 13 can be used to improve the mass estimate. Importantly, such a technique for determine accurate stellar masses can then be applied to *all* galaxies, regular or peculiar, irrespective of their dynamical state and over a range in redshift (providing the data is sufficiently precise). The technique can be considered as a modification of that frequently utilised in estimating photometric redshifts. The observed optical-infrared SED for an object of known redshift is used to optimally fit the *stellar mass* rather than the redshift in the framework of an evolutionary synthesis code. Stellar masses can be derived to within a random uncertainty of 30-50% by this technique although at present there is no reliable way to verify the results except by comparison with independent dynamical measures (Figure 14).

6. Origin of the Hubble Sequence

The availability of stellar masses for *all* types enables the construction of a powerful evolutionary plot, analogous to Figure 6, involving the *stellar mass density*, $\rho_{stars}(z, T)$, as a function of morphology T. Whilst the stellar mass density can *grow* by continued star formation, unlike the *UV luminosity density*, ρ_{UV}, it is difficult to imagine how it can *decline*. As we saw earlier ρ_{UV} can decline significantly in only 1-2 Gyr because of an abrupt truncation of activity. However, such a change would have very little effect on the infrared output as illustrated in Figure 13.

Brinchmann & Ellis (2000) secured K luminosities and optical-IR SEDS for over 300 galaxies in the CFRS/LDSS and Hawaii survey fields and derive $\rho_{stars}(z, T)$ (Figure 15). Estimating the integrated stellar mass density is prone to all of the difficulties reviewed earlier for the luminosity density and there is the added complication that the redshift surveys in question are *optically-selected* and thus must miss some (red) fraction of a true K-limited sample. Accordingly, the mass densities derived are lower limits to the true values.
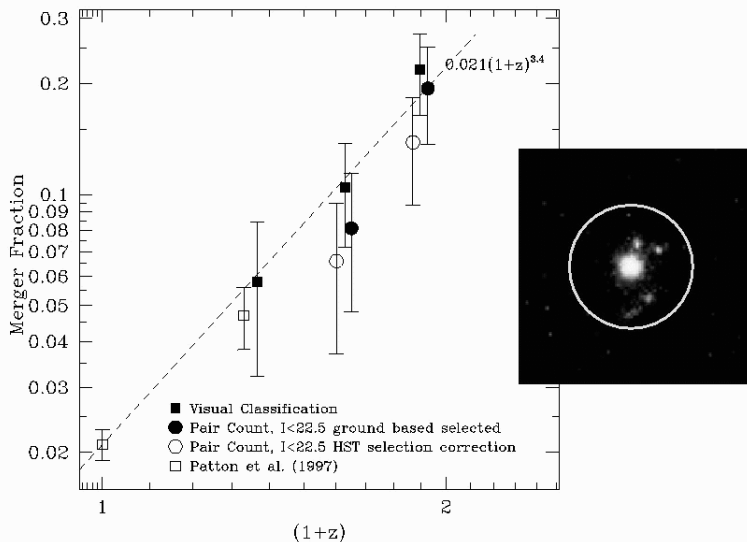

Figure 16: *An increase in the merger fraction as a function of redshift from the HST analysis of LeFevre et al. (2000). Galaxies of known redshift were examined for satellites brighter than a fixed rest-frame luminosity within a projected radius of $20h^{-1}$ kpc and corrections made for unrelated line-of-sight contamination. This redshift-dependent merger rate was adopted by Brinchmann & Ellis (2000) in Figure 15.*

Remarkably, $\rho_*(z, T)$ is a declining function for the intriguing population of morphologically-peculiar galaxies. Whereas the declining UV luminosity density could imply a fading population, such an explanation cannot be consistent with Figure 15 which argues, instead, that the objects are genuinely disappearing into other systems. The most logical explanation for their declining contribution to the stellar mass density is that morphologically-peculiar objects are being transformed, e.g. by mergers, into regular objects.

Merging has been an attractive process governing the evolution of galaxies for many years (Toomre & Toomre 1972, Rocca-Volmerange & Guiderdoni 1990, Broadhurst et al. 1992) and of course is fundamental to the hierarchical formation picture. However it has